Photophysics of Singlet and Triplet Intraligand Excited States in [ReCl(CO)3(1-(2-pyridyl)-imidazo[1,5-#]pyridine)] Complexes
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\[\text{[ReCl(CO)}_3(1-(2-	ext{pyridyl})\text{-imidazo[1,5-\alpha]pyridine})] \text{ Complexes}\]

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Abstract

Excited-state characters and dynamics of \([\text{ReCl(CO)}_3(3-R-1-(2-	ext{pyridyl})\text{-imidazo[1,5-\alpha]pyridine})] \text{ complexes (abbreviated ReGV-R, R = CH}_3, \text{ Ph, PhBu}}, \text{ PhCF}_3, \text{ PhNO}_2, \text{ PhNMe}_2)\) were investigated by pico- and nanosecond time resolved infrared spectroscopy (TRIR) and excited-state DFT and TD-DFT calculations. Near UV excitation populates the lowest singlet state \(S_1\) that undergoes picosecond intersystem crossing (ISC) to the lowest triplet \(T_1\). Both states are initially formed hot and relax with \(\sim 20\) ps lifetime. TRIR together with quantum chemical calculations reveal that \(S_1\) is predominantly a \(\pi\pi^*\) state localized at the 1-(2-pyridyl)-imidazo[1,5-\alpha]pyridine ( = impy) ligand core, with impy\(\rightarrow\)PhNO\(_2\) and PhNMe\(_2\)\(\rightarrow\)impy intraligand charge-transfer contributions in the case of ReGV-PhNO\(_2\) and ReGV-PhNMe\(_2\), respectively. \(T_1\) is predominantly \(\pi\pi^*(impy)\) in all cases. It follows that excited singlet and corresponding triplet states have to some extent different characters and structures even if originating nominally from the same preponderant one-electron excitations. ISC occurs with a solvent-independent (CH\(_2\)Cl\(_2\), MeCN) 20-30 ps lifetime, except for ReGV-PhNMe\(_2\) (10 ps in
CH$_2$Cl$_2$, 100 ps in MeCN). ISC is 200-300 times slower than in analogous complexes with low-lying MLCT states. This difference is interpreted in terms of spin-orbit interaction and characters of orbitals involved in one-electron excitations that give rise to $S_1$ and $T_1$ states. ReGV-R present a unique case of octahedral heavy-metal complexes where the $S_1$ lifetime is long enough to allow for separate spectroscopic characterization of singlet and triplet excited states. This study provides an insight into dynamics and intersystem crossing pathways of low-lying singlet and triplet excited states localized at bidentate ligands bound directly to a heavy metal atom. Rather long $^1$IL lifetimes indicate the possibility of photonic applications of singlet excited states.

Introduction

Understanding the nature and behavior of spin-singlet and triplet excited states and factors governing their conversion dynamics (intersystem crossing, ISC) in heavy-metal complexes is an intriguing scientific problem, highly relevant to developing various photonic applications. Singlet excited states are directly excited by photon absorption, serving as excitation energy gateways, whereas triplet states of these molecules often are long-lived and emissive. Their efficient population by ISC is a necessary prerequisite for using metal complexes as luminophores in OLEDs, phosphorescence imaging of biological structures, or as photosensitizers and triggers of electron transfer reactions. Excited singlet states are usually too short lived to have chemistry of their own or even to be detected, except by femtosecond fluorescence upconversion. Still, they can behave differently from corresponding triplets and could play a role in light energy harvesting by undergoing ultrafast electron transfer, competitive with ISC. Trends in ISC rates in series of structurally related metal complexes present another intriguing problem as they generally do not correlate with the spin-orbit coupling (SOC) energy. For example, ISC in Fe$^{II}$
and Ru$^{II}$ bipyridine complexes is comparable ($\leq$30 fs) and slightly faster$^{1,8,9}$ than in analogous Os$^{II}$ compounds (50-100 fs)$^{10}$ or Re$^1$ carbonyl-diimines (80-150 fs),$^{2,4}$ opposite to the SOC energy increasing in the order Fe<Ru<Os. Moreover, within the [ReX(CO)$_3$(bpy)] series, the ISC rate increases in the order $X = I$ (150 fs) < Br (130 fs) < Cl (90 fs), that is with decreasing SOC introduced by the halide ligand.$^2$ On the other hand, changing the bpy ligand for 1,10-phenanthroline or 4,7-dimethyl-1,10-phenanthroline has almost no effect.$^4$ Recent theoretical and experimental studies indicate$^{11,12}$ that the very notion of singlet and triplet states in metal diimine complexes could be inadequate, as was recognized already in 1974 by Crosby.$^{13}$ Interpreting photophysics in terms of spin-orbit states instead of singlets and triplets would allow us to discuss ISC and internal conversion in a common way and focus on the role of other effects such as vibronic coupling and distribution of excited states.$^{11,12,14}$

Rhenium(I) tricarbonyl diimine complexes are excellent systems to investigate photophysical mechanisms as well as important luminophores and sensitizers. They exhibit rich photobehavior that stems from the occurrence of different excited states whose relative energies, properties and reactivity can be controlled by variations of the structure and/or the medium.$^{15-17}$ Complexes with lowest-lying metal to ligand charge transfer (MLCT) excited states were studied most and often are used as electron-transfer triggers,$^{18-23}$ photocatalysts,$^{24,25}$ as well as luminescent probes, sensors and biomolecular imaging agents.$^{26-28}$ Re$^1$ complexes with lowest $\pi\pi^*$ intraligand (IL) excited states are less common, but important as DNA probes$^{29-31}$ and molecular photoswitches.$^{32-35}$ Low-lying excited states acquire a predominantly IL character in those [ReL(CO)$_3$(diimine)]$^{30}$ complexes where the diimine ligand consists of an extensively $\pi$-delocalized aromatic system, is a poor electron acceptor, and/or where the axial ligand L is a strong electron-acceptor, such as CO, alkyne or isonitrile.$^{17}$ Photophysics of low-lying IL states
are virtually unknown except for DNA-intercalators \([\text{Re}^1(\text{pyridine})(\text{CO})_3(R_2\text{dppz})]^+\) (dppz = dipyrido-[3,2-\(\alpha\):2',3'-\(\alpha\)]phenazine)\(^{29,30,36-38}\) whose lowest-lying \(^3\)IL states are populated alongside or via close-lying \(^3\)MLCT states. Recently synthesized \(\text{Re}^1\) carbonyl-dipyrrinate and aza dipyrromethene complexes present another interesting class of complexes with low-lying singlet and triplet IL states. They show strong visible absorption, weak red phosphorescence, and undergo CO photosubstitution upon UV (but not Vis) irradiation.\(^{39,40}\)

The complexes \([\text{ReCl(CO)}_3(3-R-1-(2-pyridyl)-imidazo[1,5-\(\alpha\)]pyridine)]\) (Figure 1, further abbreviated \(\text{ReGV-R}\)) offer a unique opportunity to systematically investigate properties and dynamics of IL states. The chelating 1-(2-pyridyl)-imidazo[1,5-\(\alpha\)]pyridine ligand core (abbreviated impy) is electron-rich, pushing both \(^1,3\)MLCT states to higher energies, above the \(^1,3\)IL states, as was demonstrated by UV-Vis absorption and photoluminescence spectroscopy, together with TD-DFT calculations.\(^{41,42}\) By changing the R-group, it is possible to append electron donating (Me), withdrawing (-Ph, -PhCF\(_3\), -PhNO\(_2\)), or reducing (-PhNMe\(_2\)) groups to the impy core (Figure 1). Using pico-nanosecond TRIR spectroscopy and excited-state DFT calculations, we have separately characterized singlet and triplet IL excited state and determined the ISC rate, which was found to be 200-300 times slower than in analogous complexes with MLCT lowest excited states.
Figure 1. Schematic structures and abbreviations of the investigated ReGV-R complexes.

Experimental

Materials

The investigated complexes (Figure 1) were synthesized and characterized using literature methods.\(^{41,42}\) Samples for TRIR measurements were prepared in MeCN and CH\(_2\)Cl\(_2\) (DCM) of spectroscopic quality (Sigma Aldrich), under air, unless stated otherwise. Fluorescence spectra were measured from solutions made of degassed MeCN (Aldrich SureSeal) in a controlled atmosphere (0.3 ppm O\(_2\)) glovebox (Jacomex).

Time-Resolved IR spectroscopy, TRIR

TRIR measurements in the v(CO) spectral region were carried out using the PIRATE instrument at the STFC Rutherford Appleton Laboratory.\(^{43-45}\) In the 0-3700 ps time domain, the
sample solution was excited (pumped) at 400 nm, using frequency-doubled pulses from a
titanium sapphire laser of ~150 fs duration (FWHM) and ca. 3 µJ energy, focused at an area ~200
µm in diameter. TRIR spectra were probed with spectrally broad IR pulses (~150 fs, width ~200
cm⁻¹) obtained by difference-frequency generation from the same titanium sapphire laser and
focused to a ~100 µm diameter in the sample. In the nanosecond range, sample excitation was
performed with 355 nm, ~0.7 ns FWHM, ~3 µJ laser pulses generated by an actively Q-switched
AOT-YVO-20QSP/MOPA Nd:Vanadate diode-pumped microlaser, that was electronically
synchronized with the femtosecond probe system with less than 400 ps jitter.⁴⁶ The 400 or 355
nm pump beams were set at the magic angle to the probe. Changes in infrared absorption at
various pump-probe time delays were recorded by normalizing the outputs from a pair of 64-
element HgCdTe (MCT) infrared linear array detectors on a shot-by-shot basis at 1 kHz. Data
were collected in pump-on/pump-off pairs to minimize the effect of long-term drift in the laser
intensity. Spectra of ReGV-Ph and ReGV-PhBu⁺ were also measured on the ULTRA
instrument⁴⁷ that uses ~50 fs pulses of ~1 µJ energy at a 10 kHz repetition rate generated by
titanium sapphire laser-based regenerative amplifier (Thales), focused to an area of ~100 µm
diameter. Probe pulses cover about 400 cm⁻¹ range. Spectra at given time delays in the 0-2000 ps
range were recorded on two 128 element HgCdTe detectors (Infrared Associates). For both
setups, the sample solutions were flowed through a 0.1 mm pathlength cell with 2 mm CaF₂
windows that was at the same time scanned-rastered across the irradiated area in two dimensions
to prevent laser heating and decomposition of the sample. FTIR spectra were measured before
and after TRIR experiments to check the sample stability. Spectral and kinetics fitting procedures
were performed using MicroCal Origin 7.1.

Photoluminescence spectra and lifetimes
Stationary emission spectra were obtained on the Fluorolog-3 instrument (model FL3-11; HORIBA Jobin Yvon). Emission decay lifetimes were measured at 560 nm on an IBH 5000 U SPC instrument equipped with a cooled Hamamatsu R3809U-50 microchannel plate photomultiplier. Samples were excited at 373 nm with an IBH NanoLED-11 diode laser (80 ps fwhm).

**DFT calculations**

Electronic structures were calculated by density functional theory (DFT) method using the Gaussian 09 and ADF2013 program packages. DFT calculations employed Perdew, Burke, Ernzerhof or M05-2X hybrid functionals. For H, C, N, and O atoms, either polarized triple-ζ basis sets 6-311g(d) were used. The Re orbitals were described with quasi-relativistic effective core pseudopotentials and corresponding optimized set of basis functions. Geometry optimizations were followed by vibrational analyses in order to characterize stationary states. The solvent was described by the polarizable continuum model (PCM): dichloromethane or acetonitrile for the ground and excited states. Low-lying excitation energies were calculated by time-dependent DFT (TD-DFT) at the optimized geometries.

Spin orbit TD-DFT calculations were done with ADF. Slater type orbitals (STO) basis sets of triple-ζ quality with two polarization functions for Re and double-ζ quality with one polarization function for remaining atoms were employed. PBE0 hybrid functional together with the scalar relativistic (SR) zero-order regular approximation (ZORA) and the COSMO model for the solvent effect corrections. Several lowest excited states have also been calculated by the perturbational approach or by the relativistic two-component zeroth-order regular approximation in TDDFT method. The difference density plots were drawn using the GaussView software.
Results and Discussion

Structures, UV-Vis and IR absorptions spectra

The investigated complexes are shown in Figure 1. Previous crystallographic studies revealed that the ligand plane is slightly tilted relative to the Re(CO)₂.equatorial plane and the phenyl group is rotated with respect to the rest of the ligand by 55-60°. Previous as well as present DFT calculations. IR spectra (Figure S1) in the region of C≡O stretching vibrations, ν(CO), are typical for fac-tricarbonyls of a Cs skeletal symmetry, consisting of a sharp band at 2021-2023 cm⁻¹ and a weaker doublet with maxima at 1913-1919 and 1888-1894 cm⁻¹, attributed to totally symmetric in-phase stretching of all three CO ligands (A'(1)), antisymmetric stretching of the two equatorial COs (A'”), and out-of-phase totally symmetric CO stretching vibration (A'(2)), respectively. This assignment is supported by the present DFT vibrational analysis that also indicates a very small admixture of the axial CO stretch to the A" mode, caused by the lack of symmetry. The lowest electronic absorption band of ReGV-R complexes occurs at about 380 nm as an intense broad structured feature (Figure 2) that consists of several intraligand (IL) and MLCT transitions (Tables S1-S3).
Time-resolved spectroscopy

Ultrafast excited-state dynamics were investigated by time-resolved IR (TRIR) spectroscopy in the ν(CO) region measured at selected time delays after 400 nm laser pulse excitation that populates an intraligand (1IL) excited state. Spectra in the nanosecond range were measured upon 355 nm irradiation into the high-energy side of the lowest electronic absorption band. TRIR spectra of the ReGV-PhR and ReGV-CH3 complexes covering the ν(CO) region are presented in Figure 3; the ReGV-PhBu' spectrum is shown in Figure S2. The high-energy part of the spectra was measured also with a high resolution of ~2 cm⁻¹ per point (Figure 4). All investigated complexes exhibit a qualitatively similar behavior that can be summarized as follows:

1. TRIR spectral evolution can be understood in terms of negative bleach bands due to depleted ground state absorption and two time-dependent overlapping transient spectral patterns. The primary pattern (1) emerges within the instrument time resolution (<1 ps) and converts on a tens-
of-picosecond time scale into the secondary pattern (2) that then slowly decays to the ground state. Each pattern consists of high-energy (HE) and low-energy (LE) spectral features which partly overlap with each other and with the bleach bands (Figures 3, S2). The presence of distinct 1HE and 2HE bands is obvious in the case of ReGV-PhNMe2 whereas 1HE appears as a shoulder for other complexes. In these cases, the 1HE and 2HE features were distinguished by analyzing the high-resolution spectra by shape fitting to multiple Lorentzians and attributed to two different transient species. Typical fits are shown in Figures S3-6. Changing the solvent (DCM, MeCN) causes small shifts of ReGV-PhNMe2 transient bands (Figures 3, 4, S3) but has virtually no effect on the TRIR spectra of ReGV-Ph and ReGV-PhBu (Figures 3, S7). None of the complexes showed any transient absorption above the A'(1) bleach (>2022 cm⁻¹), excluding population of MLCT states.
Figure 3. TRIR spectra of ReGV-PhR and ReGV-Me in DCM (except top-right) measured in the 2–3700 ps interval after ~150 fs, 400 nm excitation. Arrows indicate the temporal evolution of the spectral features: red and blue for the primary and secondary patterns, black for the bleaches. Spectral resolution: 4–5 cm⁻¹ per point. The first and last spectra are shown in red and green, respectively. The intensity decrease of the 2HE feature at later times is only apparent, caused by its shift into the A'(1) bleach region. The 1LE and 2LE features overlap strongly with each other as well with A'(2) and A" bleaches. (ReGV-PhBu’ spectra are shown in Figure S2.)
Figure 4. High-resolution TRIR spectra of selected ReGV-PhR complexes measured in the 2–3700 ps interval after 400 nm, ~150 fs excitation. Arrows indicate temporal evolution of the spectral features: red and blue arrows for the primary and secondary patterns, black for the bleaches. Resolution 1.8-1.9 cm\(^{-1}\) per point. TRIR spectra of ReGV-PhBu', -PhCH\(_3\), and -PhNO\(_2\) are very similar to those of ReGV-PhCF\(_3\).

2. The primary TRIR pattern consists of two broad features between 1860-1920 cm\(^{-1}\) (1LE) and in the 1995-2011 cm\(^{-1}\) range (1HE). The 1HE feature is best developed for ReGV-PhNMe\(_2\), where it appears as a broad distinct band at 1998 cm\(^{-1}\) in MeCN (2000 cm\(^{-1}\) in DCM), shifted ~20 cm\(^{-1}\) below the ground-state A'(1) band. For all other complexes, 1HE occurs as a broad shoulder that overlaps with the 2HE band of the secondary spectral pattern. Lorentzian shape fitting (Figures S3-6) places the 1HE maximum at 2006-2008 cm\(^{-1}\) in DCM, regardless of the particular ligand substituent.

3. The secondary spectral pattern shows a broad 2LE absorption in the 1880-1920 cm\(^{-1}\) range and a sharp 2HE band at 2014-2019 cm\(^{-1}\), depending on the particular complex. (For ReGV-
PhNMe2: 2014-2015 cm$^{-1}$ in both DCM and MeCN, measured at time delays >200 ps.) The 2HE maximum typically lies 4-5 cm$^{-1}$ below the A'(1) bleach for all ReGV-PhR and 6-7 cm$^{-1}$ for ReGV-Me. Neither the shape nor position of the 2LE feature can be determined because of an extensive overlap with 1LE and the ground-state bleach. The 2LE feature lies at slightly higher wavenumbers than 1LE for all complexes except ReGV-PhNO2, where 2LE lies slightly below than (or comparable with) 1LE.

4. Temporal evolution of the IR spectra combines conversion of the primary to the secondary pattern with small dynamic shifts of all bands to higher wavenumbers and their narrowing. Whereas the pattern conversion manifests a structural change of the photogenerated transient species, the continuous band shift/narrowing is due\textsuperscript{45,63,64} to vibrational cooling and solvent relaxation. Kinetics analysis is complicated by similar timescales of all these processes and overlap between different spectral features. (For example, the bleach recovery and the 2HE decay on the ps timescale are only apparent, caused by the positive transient spectral features moving more into the bleach regions of negative absorptions, instead of transient decay to the ground state.)

The spectral pattern conversion is best observable in the case of ReGV-PhNMe2, where the IR bands are narrower and well separated. The conversion is manifested by an isosbestic 1HE decay accompanied by a 2HE rise. The 1HE and 2HE band areas decrease and increase, respectively, with a common lifetime of ~10 ps in DCM (1HE: 9.1±0.7 ps; 2HE: 10.0±0.4 ps) and ~100 ps in MeCN (1HE: 92±5; 2HE: 102±6 ps) (Figure 5). The 1LE band decays concomitantly with the 1HE decay but the kinetics analysis is affected by an overlap with shifting 2LE band: 87±2 ps (MeCN); 13-16 ps (DCM). Other ReGV-PhR and ReGV-CH3 complexes show a common behavior: A concomitant decay of the 1LE and 1HE features and 2HE rise with
lifetimes in the 20-30 ps range, measured in DCM. Changing the solvent to MeCN was examined for ReGV-Ph and ReGV-PhBu' and found to have virtually no effect on the conversion kinetics (Figure S7). (Somewhat exceptional behavior was observed for ReGV-PhNO2, where the 1LE/2LE feature is partly formed in <1 ps, then increases further and decays with lifetimes of 16±8 ps and 31±10 ps, respectively. Both 1HE and 2HE features are very broad at short time delays and overlap extensively with each other, complicating kinetics analysis. 2HE narrowing occurs with a 37±2 ps lifetime.)

Dynamic 1HE and 2HE band shifts and narrowing were quantified for ReGV-PhNMe2 using Lorentzian fits (Figure 6). The 1HE band in MeCN undergoes ~2 cm⁻¹ upshift in 15 ps and ~35% narrowing completed in the first 10 ps. The 2HE band shifts by +0.8 cm⁻¹ with a 21±1 ps lifetime. In DCM, 2HE shifts by 1.5 cm⁻¹, with a 20±1 ps lifetime and narrows by ~15% with a 30-40 ps lifetime. Similar 2HE shifts and narrowing were observed for other complexes, although the analysis is prevented by more extensive spectral overlaps.

**Figure 5.** Time-dependence of 1HE and 2HE band areas of ReGV-PhNMe2 in MeCN (left) and DCM (right). Areas were determined by Lorentzian fitting of high-resolution TRIR spectra, Figure S3.
Figure 6. Time-dependence of the maximum wavenumber (top, middle) and the width (bottom) of the 2HE band of ReGV-PhNMMe2. Band maxima and widths (fwhm) were determined by Lorentzian shape-fitting (Figure S3) and their time dependences were fitted as single-exponential functions. Lifetimes and magnitudes of the band shift/narrowing are shown in the insets.

5. On the ns time scale (Figure 7), the 2HE and 2LE features and the bleach bands decay with a common lifetime: 250±17 (ReGV-PhBu', DCM), 143±2 (ReGV-PhBu', MeCN), 270±30 (ReGV-PhCF3, DCM), 144±2 (ReGV-Ph, MeCN), and 153±15 ns (ReGV-PhNMMe2, DCM), measured in aerated solutions. Much faster decay of 7±0.6 ns was observed for ReGV-PhNO2 in DCM. Nanosecond TRIR spectra of ReGV-PhBu' in CD3CN measured in the fingerprint region (Figure S8) show a series of bleach- and down-shifted transient bands due to GV-PhBu' ligand vibrations. All these spectral features decay together without any shifts, showing that the secondary transient does not undergo any nanosecond structural changes. The exceptionally fast decay observed for ReGV-PhNO2 is typical for metal complexes with nitro-substituted ligands, presumably due to strong coupling between the –NO2 group and solvent vibrational modes.65
Figure 7. Nanosecond TRIR spectra of selected ReGV-R complexes measured in DCM at selected time delays after ~0.7 ns, 355 nm excitation. Resolution: 4-5 cm\(^{-1}\).

6. All the investigated complexes except ReGV-PhNMe2 were reported to show structured long-lived phosphorescence in deaerated solutions.\(^{41,42}\) Herein, we have revisited the photoluminescence of ReGV-Ph in MeCN (Figure 1) and determined the vibrational progression of 1376 cm\(^{-1}\). The emission intensity in air-saturated solution decays with a 157±10 ns lifetime, which is comparable with the TRIR 2HE lifetime of 144±2 ns. Comparison with the lifetime measured in a degassed solution (29-37 µs) allows us to estimate the bimolecular rate constant of oxygen quenching as 3×10\(^9\) M\(^{-1}\)s\(^{-1}\) (assuming the \(O_2\) concentration in air-saturated MeCN of 2.42×10\(^{-3}\) M). The excitation spectrum matches the absorption spectrum (Figure 2), confirming that the emitting state is populated with equal efficiency at both excitation wavelengths used in the TRIR experiments, 355 and 400 nm.
Nature of the TRIR transients

All investigated ReGV-R complexes show common photophysical behavior, whereby optical excitation produces a picosecond-lived transient species that undergoes conversion into the secondary transient ultimately decaying into the ground state. The secondary species can be identified with the lowest triplet state T1 because of its emission, long lifetime, fast oxygen quenching, and direct decay to the ground state. Conversion of the primary to the secondary species is clearly manifested in TRIR spectra by the 1HE decay accompanied by the 2HE rise and by similar changes in the LE region. It can correspond either to a conformational change (restructuring) of the T1 state or to a nonradiative transition from a higher excited state. Given the rather rigid molecular structures of the ReCl(CO)3(impy) core of the investigated complexes, conformational changes can involve rotation of the -PhR group relative to the impy plane, rotation of the -NO2 group in ReGV-PhNO2, or -NMe2 rotation in ReGV-PhNMe2. However, involvement of such rotational motions is excluded by the common spectral and dynamic characteristics of the primary and secondary spectral patterns observed for all investigated complexes, including those with a non-planar substituent R (ReGV-Ph, ReGV-PhBu\textsuperscript{1}, ReGV-PhCF3) and ReGV-CH3 with no pendant aromatic group. These considerations thus leave a higher excited state as the most likely assignment of the primary transient.

All ν(CO) features in both TRIR spectral patterns are shifted to lower wavenumbers from the corresponding ground-state values indicating increased Re→CO π back donation and excluding any involvement of MLCT states. Such downshifts suggest either IL\textsuperscript{36-38} sigma-bond-to-ligand charge transfer (SBLCT),\textsuperscript{67} or ligand-to-ligand charge transfer (LLCT)\textsuperscript{68} excited states, or a ligand-localized reduction.\textsuperscript{18-20,38,69} SBLCT and LLCT states do not occur at low energies in
rhenium carbonyl diimines with chloride as the axial ligand.\textsuperscript{17,70} GV-R ligand reduction is out of question because of very negative redox potentials (\(<-2\ V\ vs.\ Fc^+/Fc\)),\textsuperscript{41} and an absence of a reducing group, with the exception of ReGV-PhNMe\textsubscript{2} that, however, behaves similarly to complexes containing electron-withdrawing (-PhCF\textsubscript{3}) or oxidizing (-PhNO\textsubscript{2}) substituents. Hence, both the primary and secondary excited states can be assigned as IL states. A slightly larger ν(CO) down-shift observed for the primary IR pattern of ReGV-PhNMe\textsubscript{2} indicates a more extensive intra-ligand electron density redistribution towards the impy part of the ligand than in other ReGV-R complexes, likely due to the electron-donating character of the NMe\textsubscript{2} group. For the secondary state (T\textsubscript{1}), the IL assignment is further supported by the characteristic structured emission and microsecond lifetimes that are much longer than those\textsuperscript{12,17,71} of analogous complexes with a \(^3\)MLCT lowest excited state (usually tens of nanoseconds).

More detailed characterization of excited states involved in ReGV-R photophysics emerges from excited-state DFT and TD-DFT calculations that were performed on ReGV-Ph, ReGV-PhNO\textsubscript{2}, and ReGV-PhNMe\textsubscript{2}. Electron density redistribution upon the lowest transition to the S\textsubscript{1} singlet excited state is shown in Figure 8 and summarized in Table S4. It follows that the S\textsubscript{1} transition involves π→π* excitation of the impy part of the GV ligand that is combined with a PhNMe\textsubscript{2}→impy and impy→PhNO\textsubscript{2} intraligand charge transfer in the case of ReGV-PhNMe\textsubscript{2} and ReGV-PhNO\textsubscript{2}, respectively. Electron density on N-donor atoms and (less) on equatorial CO ligands increases in each case, accompanied by a small electron depopulation of a Re 5dπ orbital, reflecting increasing Re→CO π back donation.
Figure 8. Electron density differences between the $S_1$ excited states of ReGV-Ph, ReGV-PhNMe2, and ReGV-PhNO2 and the corresponding ground states. Regions of depleted and increased electron density are shown in blue and red, respectively. (TD-DFT calculation at optimized ground-state geometries, M052x, 6-31g(d), PCM for CH$_2$Cl$_2$)

Excited-state energy diagrams (Figure 9) indicate the electronic states that could participate in the observed photophysics. At the ground state geometry, there are several triplet states lying between the optically populated $S_1$ and the ground state: $T_1$ and $T_2$ for ReGV-Ph, $T_1$-$T_3$ for ReGV-PhNMe2, and $T_1$-$T_5$ for ReGV-PhNO2. The $S_1$ state is strongly stabilized in energy upon relaxation. Structurally optimized $S_1$ is separated from the ground state only by the lowest triplet $T_1$ in the cases of ReGV-Ph and ReGV-PhNMe2, and by $T_1$, $T_2$, and $T_3$ for ReGV-PhNO2, suggesting the $S_1$$\rightarrow$$T_1$ intersystem crossing (ISC) to be the most likely mechanism responsible for the conversion observed in the TRIR spectra.

This interpretation is fully supported by the TD-DFT characterization of the $S_1$ and $T_1$ excited states whose calculated $\nu$(CO) IR spectra match well the experimental primary and secondary $\nu$(CO) spectral patterns, respectively. Figures 10 and S10 show that calculated IR spectra of the structurally optimized $S_1$ and $T_1$ states of ReGV-PhNMe2 and ReGV-Ph reproduce the conversion of the 1HE band to the higher-lying 2HE band and the small shift of the LE features to higher wavenumbers upon the $S_1$$\rightarrow$$T_1$ ISC. The match between the calculated and experimental $S_1$ spectra of ReGV-PhNO2 (Figure S11) is not as good as for the other two
complexes, likely due to extensive relaxation-induced shifts and/or involvement of an intermediate state (T$_2$ or T$_3$) in the ISC that affect experimental spectra. The occurrence of such processes was indicated experimentally by the rise and decay of the 1LE+2HE feature and the large ~30 ps dynamic narrowing of both HE features, which are discussed above. In fact, the state diagram (Figure 9-right) supports the possibility that the ISC in this complex could occur via T$_2$ and T$_3$ higher excited states. Still, the calculated ReGV-PhNO$_2$ spectra correctly predict the T$_1$ 2LE band occurring at slightly lower wavenumbers than S$_1$ 1LE feature, in contrast to other complexes.

Figure 9. State energy diagram of ReGV-Ph, ReGV-PhNMe$_2$, and ReGV-PhNO$_2$ in CH$_2$Cl$_2$. Electronic states calculated at optimized ground-state, S$_1$, and T$_1$ geometries are shown for each complex. Black: ground state; blue: triplet states; red: excited singlet states. (TD-DFT, M052x, 6-31g(d), PCM for CH$_2$Cl$_2$)
Figure 10. Calculated and experimental excited-state IR spectra of ReGV-PhNMe2 in MeCN. Top: calculated spectra of S\textsubscript{1} at the optimized S\textsubscript{1} geometry (red) and of T\textsubscript{1} at the optimized T\textsubscript{1} geometry (blue). Middle: calculated difference IR spectra of S\textsubscript{1} (red) and T\textsubscript{1} (blue) minus the ground state spectrum calculated at the optimized ground-state geometry. Bottom: experimental TRIR spectra measured at 2 ps (red) and 3000 ps (blue). The red spectrum corresponds to a hot state and therefore is slightly broadened and shifted to lower wavenumbers. The calculated and experimental 1HE and 2HE wavenumbers match with an accuracy of 4 and 2 cm\textsuperscript{-1}, respectively: The 1HE band (S\textsubscript{1} state) was calculated at 2002 cm\textsuperscript{-1} vs. the experimental value of 1998 cm\textsuperscript{-1} (determined by shape-fitting of high-resolution spectra). For 2HE, the calculated and experimental values are 2012 and 2014 cm\textsuperscript{-1}, respectively. Calculated and experimental 1LE wavenumbers are 1878 and 1874 cm\textsuperscript{-1}, respectively. Calculation: TD-DFT, 6-31g(d), M052X, PCM for MeCN, scaling factor 0.940.

The relaxed S\textsubscript{1} state of ReGV-Ph and ReGV-PhNMe2 is calculated to have a predominantly impy-localized \pi\pi* intraligand character combined with a small increase of Re\rightarrow CO \pi donation (Table S4). This conclusion can be extended to all other investigated complexes, except ReGV-PhNO\textsubscript{2} whose relaxed S\textsubscript{1} state also is predominantly intraligand impy-localized \pi\pi*, but accompanied by a limited Re\rightarrow GV-PhNO\textsubscript{2} charge transfer (MLCT), leaving the electron density on CO ligands intact. The T\textsubscript{1} state of all three complexes was
calculated to be an impy-localized \( \pi\pi^* \) \( ^3 \)IL, (Figure 9, Table S5). This conclusion is confirmed for relaxed T\(_1\) by spin-density distributions calculated at optimized T\(_1\) geometries (Figure S9, Table S5). The larger involvement of the Re(CO)\(_3\) moiety in S\(_1\) than T\(_1\) accounts for the fact that the S\(_1\) \( \nu(\text{CO}) \) bands of the primary TRIR pattern occur at lower wavenumbers those of the secondary pattern due to T\(_1\). The situation is different only in the case of ReGV-PhNO\(_2\), where the Re\(\rightarrow\)GV-PhNO\(_2\) CT contribution, present in S\(_1\) vanishes upon ISC to T\(_1\). Consequently, the S\(_1\) 1LE band occurs at higher wavenumbers than the T\(_1\) 2LE feature (Figure S11).

\[ \text{Figure 11. Electron density differences between the three lowest triplet states of ReGV-Ph, ReGV-PhNMe2, and ReGV-PhNO}_2\text{ and the corresponding ground states. Regions of depleted and increased electron density are shown in blue and red, respectively. (TD-DFT calculation at optimized ground-state geometries, M052x, 6-31g(d), PCM for CH}_2\text{Cl}_2) } \]

Finally, it is possible to conclude that any involvement of the T\(_2\) state in the ISC is very unlikely in all investigated complexes except ReGV-PhNO\(_2\). The T\(_2\) electron density difference
maps of \( \text{ReGV-Ph} \) and \( \text{ReGV-PhNMe}_2 \) (Figure 11) do not show any electron density increase on CO ligands while indicating MLCT \( \text{Re} \rightarrow \text{GV} \) contributions (manifested by Re d-orbital depopulation and electron density increase on the ligand). This type of electron density distribution qualitatively predicts that the \( T_2 \) \( \nu(\text{CO}) \) bands would lie at higher wavenumbers than those of \( T_1 \) and cannot thus account for the primary TRIR pattern. Moreover, whereas \( T_1 \) is in all cases an impy-localized \( \pi\pi^* \) state, the intraligand charge distribution in higher triplets is expected to depend on the ligand substituent. However, the observation of a similar TRIR behavior was observed for complexes with oxidizable (\(-\text{PhNMe}_2\)), reducible (\(-\text{PhNO}_2\)), electron-withdrawing (\(-\text{PhCF}_3\)), donating (\(-\text{PhBu}^t\)), and "inert" (\(-\text{Ph}, \ -\text{Me}\)) groups effectively excludes the conversion between two triplet excited states of different charge distributions as mechanism responsible for TRIR pattern conversion.

**Spin-orbit effects**

Perturbational and two-component spin-orbit DFT calculations of \( \text{ReGV-PhNMe}_2 \) were carried out using pure (BP86) as well as hybrid (B3LYP) functionals in order to further examine the validity of the above interpretation and to assess SO effects on intraligand states. In all SO calculations performed, the \( S_1 \) and \( T_1 \) states identified above remain the two lowest excited states. Zero field splitting (zfs) of the \( T_1 \) state was calculated to be very small (~0.216 cm\(^{-1}\)) and comparable with experimentally determined zfs of \(^3\text{IL}\) states of \( \text{Rh}^{\text{III}} \) bpy and phen complexes (0.1-0.2 cm\(^{-1}\))\(^{72,73}\) or \([\text{Ir}(2\text{-phenylpyridine})_2(\text{CO})(\text{Cl})]\) (< 1 cm\(^{-1}\))\(^{74}\) In \( \text{ReGV-PhNMe}_2 \), zfs arises from a very small (~0.2%) SO-induced admixture of higher-lying \(^1\text{MLCT}\) singlet state(s), instead of the optically populated \( S_1 \). The \( T_1 \) zfs is much smaller than that calculated (133 cm\(^{-1}\))\(^{12}\) or measured (90 cm\(^{-1}\))\(^{75}\) for the predominantly MLCT lowest triplet state of \([\text{ReCl}(\text{CO})_3(\text{bpy})]\), which originates from mixing with the optically excited \(^1\text{MLCT}\) state. The \( S_1 \) state of \( \text{ReGV-} \)
PhNMe2 was calculated to be essentially a pure singlet using the BP86 functional that, however, strongly underestimates the transition energy. B3LYP SO calculations indicate a 23% admixture of higher triplets to S\textsubscript{1} and occurrence of three more SO states with predominant triplet characters at only slightly higher energies (240-400 cm\textsuperscript{-1}). The sensitivity of S\textsubscript{1} SO calculations to the functional precludes making any detailed conclusions on the S\textsubscript{1} spin character. However, a significant admixture of higher MLCT triplets into S\textsubscript{1} is very unlikely in view of the IR spectral pattern that is typical of an IL state and the match of the experimental spectrum with that calculated for spin-free S\textsubscript{1} (Figure 10).

\textit{Ultrafast excited-state dynamics}

TRIR experiments reveal two kinds of picosecond dynamics of ReGV-R complexes: (i) decay of the S\textsubscript{1} state into the lowest triplet T\textsubscript{1} manifested by the conversion of the primary TRIR spectral pattern to the secondary one, and (ii) structural/solvational relaxation of both S\textsubscript{1} and T\textsubscript{1} that is revealed by small dynamic 1HE and 2HE band shifts and narrowing.

Relaxation-induced spectral shifts occur on a timescale of ~20 ps that is slightly slower than those observed\textsuperscript{45,63,64,76,77} in MLCT excited-state TRIR spectra of Re carbonyl-diimines in dipolar solvents (9-15 ps). The shifts are attributed\textsuperscript{45,76,77} to cooling that includes restructuring of the first solvent shell to optimize interactions between the solvent molecules and electronically excited ReGV-R solute whose electron distribution over the ligand is different from that in the ground state.

The S\textsubscript{1} to T\textsubscript{1} conversion (intersystem crossing, ISC) occurs with a 20-30 ps lifetime for ReGV-CH\textsubscript{3}, -Ph, -PhBu\textsuperscript{t}, -PhCF\textsubscript{3}, and –PhNO\textsubscript{2} in DCM, and about 10 ps in the case of ReGV-PhNMe2. Changing the solvent to more polar MeCN has no effect on the ISC rate of
ReGV-Ph and ReGV-PhBu′ but slows down ISC in ReGV-PhNMe2 10-times to ~100 ps. ISC in ReGV-R is 200-300 times slower than in [ReCl(CO)3(bpy)] and [ReCl(CO)3(phen)] (85-110 fs, depending on the solvent). This huge change in the ISC rate is attributable to two effects: different excited-state characters and different symmetry constraints. The femtosecond ISC in bpy and phen complexes involves singlet and triplet states of a predominantly MLCT character with ~50% Re 5d orbital participation in the depopulated molecular orbitals. More importantly, the optically populated 1MLCT and the lowest 3MLCT states differ in the symmetry of the 5d orbital involved and, hence, ISC is accompanied by orbital rotation that compensates for the change in the spin momentum. The 1,3MLCT states are thus directly SO-coupled and the ISC is allowed. The situation is very different in the case of ReGV-R complexes, where both S1 and T1 are of an IL character, arising predominantly from HOMO→LUMO excitation (Tables S1-3). The molecular orbitals depopulated in S1 and T1 are the same, no orbital rotation occurs and ISC is forbidden according to El Sayed's rules in the same way as 1ππ*→3ππ* ISC in aromatic hydrocarbons. Moreover, the Re 5d participation in S1 and T1 is low, originating mostly from small interconfigurational mixing between HOMO→LUMO and higher MLCT excitations. In the absence of direct S1-T1 SO interaction (confirmed by SO calculations), ISC is probably enabled by weaker second-order interactions of S1 with higher-lying MLCT states. The difference in the SO coupling strength and, hence, the propensity to undergo ISC between the singlet and triplet IL states in ReGV-R and MLCT states in [ReCl(CO)3(bpy)] also is demonstrated by the vastly different zfs magnitudes, 0.2 vs. 133 cm−1, see above.

The S1 and T1 states in ReGV-Ph, -PhBu′, and -PhCF3 behave as ππ* excited states with small Re(CO)3 contributions. There are no large differences in the charge distribution between the ground states and the two excited states in question (Tables S4, S5). Consequently, the ISC
rates of ReGV-Ph and ReGV-PhBu\(^t\) are insensitive to the solvent polarity, DCM vs. MeCN. On the other hand, the pronounced solvent effect on the ISC rate in ReGV-PhNMe\(_2\) likely is due to the larger intraligand charge-transfer character of S\(_1\) than T\(_1\) that makes their energy gap and vibronic coupling sensitive to the solvent polarity.

**Concluding remarks**

ReGV-R are an unusual class of heavy metal complexes in which the lowest lying singlet as well as triplet excited states are predominantly localized at a N,N-chelating ligand that is directly coordinated to the metal atom. In this respect, ReGV-R differ from the other prominent class of compounds showing ligand-localized photophysics and photochemistry, namely [Pt(L\(_2\))(C≡CR')\(_2\)] and [Pt(terpyridine)(C≡CR')]\(^+\), where the \(\pi\pi^*\) IL states are localized predominantly at a pendant aromatic group R', such as naphthalene\(^79\), pyrene\(^80\), perylenediimide,\(^81,82\) or stilbene\(^83\) (L\(_2\) = 4,4'-Bu\(^t\)-bpy or diphosphine). ReGV-R photophysics also differ from those of DNA intercallators [Re(L)(CO)\(_3\)(dppz)]\(^+\) whose lowest \(^3\)IL state is populated on a sub-picosecond timescale via MLCT states and, in some cases, occurs in a thermal equilibrium with \(^3\)MLCT.\(^29,30,36-38\)

ReGV-R present a rare case where both the optically populated excited singlet and the lowest triplet states were characterized spectroscopically. It turns out that these two states differ considerably in their respective characters, despite being nominally derived from the same preponderant one-electron excitation, \(i.e.\) HOMO→LUMO. This difference is most pronounced for ReGV-NMe\(_2\), where both S\(_1\) and T\(_1\) are essentially \(\pi\pi^*\) excited states of the impy ligand core but S\(_1\) contains an additional PhNMe\(_2\)→impy charge transfer component. The Re(CO)\(_3\) electronic involvement, although small in absolute terms, is larger for S\(_1\) than T\(_1\). Such differences result
from different extent of interconfigurational mixing of higher one-electron excitations into $S_1$ and $T_1$. On the other hand, ReGV-R complexes without strongly donating (or accepting) substituents such as ReGV-Ph or ReGV-CH$_3$ have both $S_1$ and $T_1$ states localized mostly at the impy part of the ligand but still with somewhat different IR spectra and, hence, electronic structures. Quadruply bonded W$^{II}$ dimers with electron-accepting bridging ligands present another case showing differences in singlet and triplet excited-state IR spectra that indicate a larger electron density on the ligand in $1\delta\pi^*$ than $3\delta\pi^*$ MLCT state.$^{84,85}$ In the case of 2,5-bis(p-arylethynyl)cyclopentadiene complexes of Rh$^{III}$, a larger $\nu$(C≡C) IR shift was observed on going from the ground state to the lowest $3^1$IL state than to the corresponding excited singlet.$^{86}$

The slow ISC between the $S_1$ and $T_1$ IL states in ReGV-R, as compared with MLCT states in [ReCl(CO)$_3$(bpy)] and similar complexes, is due not only to smaller Re 5d participation but mainly to different symmetry constraints which allow for a direct SO coupling between the optically populated $1^1$MLCT and the lowest $3^1$MLCT states but not between IL singlets and triplets. The relative orientation (rotation) of the depopulated orbital in the singlet and triplet is the decisive factor. Direct SO coupling between $1^1,3^1$MLCT states leads to a large zfs and femtosecond ISC that can be treated as a vibronic transition between two spin-orbit states.$^{11,12}$ On the other hand, the intraligand $T_1$ zfs is very small and the $S_1\rightarrow T_1$ ISC involves second-order coupling via higher excited states. Indeed, an identical preponderant electron configuration of the $S_1$ and $T_1$ states seems to be the key factor extending excited singlet lifetimes in Re-GV as well as in other types of complexes and electronic states, for example: $1^1$MLCT of pseudotetrahedral complexes of Cu$^1$ (13-16 ps)$^{87}$ and Pt$^0$ (3.2 ps)$^{88}$, $1^1d\sigma^*p\sigma$ in $d^8$-$d^6$ dimers such as [Pt$_2$(P$_2$O$_5$H$_2$)$_4$]$^{4+}$ (3-30 ps) and [Pt$_2$(P$_2$O$_5$BF$_2$)$_4$]$^{4+}$ (1.6 ns),$^{89}$ and for IL states in Pt complexes with pendant organic chromophores: [Pt(PBu$_3$)$_2$(C≡C-pyrene)$_2$] (5.4 ps),$^{80}$ [Pt(L$_2$)(C≡C-peryleneimide)$_2$] (2-4 ps),$^{81}$
or [Pt(trpy)(C≡C-perylenediimide)]⁺ (109 ps).² Surprisingly, \(^1\text{IL} \rightarrow ^3\text{IL ISC}\) in ReGV-R is slower than in some of the Pt complexes, despite of a direct linkage of the photophysically active ligand to the Re atom.

Metal complexes with lowest triplets of a \(\pi\pi^*\) intraligand character are deemed not to be good candidates for OLED luminophores because of low phosphorescence quantum yields. This is a consequence of weak SO coupling and negligible singlet admixtures into the lowest triplet state that also are manifested by very small (< 1 cm⁻¹) zfs.⁹⁰-⁹³ The present study of ReGV-R photophysics confirms this picture, provides a deeper insight into spin-orbit coupling pathways in metal complexes with low-lying intraligand states, and reveals rather long \(^1\text{IL}\) lifetimes. Metal complexes with picosecond-lived optically populated singlet excited states could find interesting photonic applications. For example, they could be considered as active components of molecular antennas or artificial reaction centers where the absorbed light energy would be efficiently captured and utilized either by energy- or electron transfer before being partly lost by conversion to lower-lying triplets. Long singlet lifetimes also could allow for singlet fission⁹⁴ that has yet to be observed in metal complexes.

**Supporting Information.** Tables of DFT-calculated low-lying singlet-singlet and singlet transitions (energies, osc. strengths), table of Mulliken charges in the \(S_1, T_1\), and the ground-state, figures of ground-state IR spectra, TRIR of ReGV-PhBu in DCM, typical Lorentzian shape-fits of TRIR spectra, TRIR spectra of ReGV-Ph and ReGV-PhBu in MeCN, TRIR spectra of ReGV-PhBu in CD₃CN in the fingerprint region, calculated spin density distribution in relaxed \(T_1\) excited states, comparison of calculated and experimental excited-state IR spectra of ReGV-Ph and ReGV-PhNO₂. Full reference for the Gaussian software. This information is available free of charge via the internet at http://pubs.acs.org.

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Graphical Abstract

100 ps ISC between intraligand excited states